The Magnetic Field Strength and Energy Balance of OMC 1



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<u>BISTRO:</u> Derek Ward-Thompson, Woojin Kwon, Tetsuo Hasegawa, Shih-Ping Lai, Keping Qiu, Pierre Bastien, Ray Furuya, Simon Coudé, Jongsoo Kim, Chang Won Lee, Andy Pon, Sarah Sadavoy, among many others

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<u>Gould Belt Survey:</u> James Di Francesco, Jenny Hatchell, Helen Kirk, among many others

BISTRO: Overview



- We are a JCMT Large Program aimed at mapping the Gould Belt star-forming regions in polarised light
- >100 survey members across 6 partner regions and the East Asian Observatory
- P.I.s: Derek Ward-Thompson (UK), Tetsuo Hasegawa (Japan), Woojin Kwon (Korea), Shih-Ping Lai (Taiwan), Keping Qiu (China), Pierre Bastien (Canada)
- BISTRO-1 awarded 224 hours of observing time to map: Ophiuchus, Orion A & B, Perseus, Serpens Main, Taurus L1495/B211, Auriga, IC5146
- BISTRO-2 (2017) awarded a further 224 hours to map: More of Orion, more of Perseus, Serpens Aquila, M16, DR15, DR21, NGC 2264, NGC 6334, Mon R2, Rosette

Survey paper: Ward-Thompson et al. 2017, ApJ 842 66

BISTRO: Scientific Goals



- To map the magnetic field within cores and filaments, on scales of ~1000-5000 AU
- To determine magnetic field strengths in nearby molecular clouds using the Chandrasekhar-Fermi method (through synthesis with Gould Belt Survey HARP data)
- To investigate the relative importance of magnetic fields and turbulence to star formation
- To test the model of magnetic funnelling of material onto filaments (André et al. 2013; Palmeirim et al. 2013)
- To investigate the role of magnetic fields in shaping protostellar evolution
- To investigate the effect of magnetic fields on bipolar outflows from young protostars

POL-2: The Instrument



- A single-beam imaging polarimeter mounted on the SCUBA-2 camera on the JCMT (15m)
- Measures linear polarisation (Stokes Q & U)
- Takes data at 850μm (353 GHz) and 450μm (667 GHz) simultaneously. 850μm commissioned; 450μm commissioning ongoing
- 14" resolution at 850µm, sensitive to spatial scales ≲5', mapping mode produces 12'-diameter maps with the option of mosaicing
- For details: Friberg et al. 2016, SPIE 9914 03

Orion A and OMC 1

The Orion Nebula is the nearest high-mass star-forming region.

OMC 1 is a dense molecular cloud which is a site of ongoing high-mass star formation located in the centre of the Orion A "integral filament", at a distance of 388 pc (Kounkel et al. 2017)

OMC 1 is located behind the Trapezium cluster, and bounded on one side by the Orion Bar PDR.

Image: CO-subtracted 850µm SCUBA-2 data; JCMT GBS IR2 reduction



The magnetic field of OMC 1



- Observed in January 2016 as part of POL-2 commissioning and BISTRO science programmes
- ~2mJy/beam RMS sensitivity on 12" pixels
- Vectors rotated to trace magnetic field
- Note 'hourglass' morphology (c.f. Schleuning 1998)

The BN/KL outflow

- a wide-angle explosive outflow with multiple ejecta (the "bullets of Orion"; Allen & Burton 1993)

 one of the most energetic outflows observed in a starforming region: contains ~4x10⁴⁷ erg of energy (Kwan & Scoville 1976)

likely to have been formed by an encounter between Orion sources BN, n and I around 500 years ago (Gomez et al 2005; Bally & Zinnecker 2005)



Hypothesis 1: Outflow shapes field



Hypothesis 2: Field shapes outflow a) b) C) Gravitationally Unstable **Dense Filament** Magnetic Field **Explosive** Overdense Outflow Clumps

Chandrasekhar-Fermi Method

Assumes equipartition between non-thermal motions and the magnetic field: deviation in angle from the mean field direction is taken to be the result of distortion of the field by small-scale non-thermal motions (see Chandrasekhar & Fermi 1953).





Deviation in angle: Unsharp masking approach



Mean field direction determined by smoothing with a 3x3 pixel box.

Pixels with >90° changes in angle inside the smoothing box excluded.

Deviation in angle: Unsharp masking approach





$$\Delta \theta = \theta_{obs} - \langle \theta \rangle$$

In pixels where P/DP > 5 and $\delta \theta_{\text{max}} < 2.0^{\circ}$:

 $\langle \sigma(\Delta \theta) \rangle$: 4.0° ± 0.3°



Magnetic field strength estimates

$$B_{\rm pos} = Q \sqrt{4\pi\rho} \, \frac{\delta v}{\delta\theta} \approx 9.3 \sqrt{n({\rm H}_2)} \, \frac{\Delta v}{\langle \Delta \theta \rangle}$$

 $\Delta v = 3.12 \pm 0.73 \text{ km s}^{-1}$

 $\langle \sigma(\Delta \theta) \rangle = 4.0^{\circ} \pm 0.3^{\circ}$

 $n(H_2) \approx (0.8 \pm 0.6) \times 10^6 \text{ cm}^{-3}$

Best estimate: $B_{pos} = 6.6 \pm 4.7 \text{ mG}$

This value is consistent with previous estimates for OMC 1 (Hildebrand et al. 2009), and with values measured in other star-forming regions (Curran & Chrysostomou 2007)

Magnetic energy density

Using our measured magnetic field strength, we can estimate magnetic energy density in OMC 1:

$$U_B = \frac{|\mathbf{B}|^2}{2\mu_0}$$

If B = 6.6 mG, then $U_{B} = 1.7 \times 10^{-7} \text{ J m}^{-3}$

How does this compare to the other sources of energy in OMC 1?

Gravitational potential energy density of the BN/KL – S system

From our column density map: Mass of BN/KL: 1000 M_{\odot} Mass of S: 285 M_{\odot}

Treating the system as a uniform cylinder: $U_{G} = 1.9 \times 10^{-7} \text{ J m}^{-3}$

Treating the system as a pair of point sources: $U_{G} = 0.5 \times 10^{-7} \, J \, m^{-3}$

Treating the system as a prolate spheroid: $U_G = 8.8 \times 10^{-7} \text{ J m}^{-3}$



Outflow energy density

$$\sim 4\! imes\!10^{40}\,\mathrm{J}\,$$
 (Kwan & Scoville 1976) $_{_2}$

$$V_{\text{outflow}} = 2 \times \frac{2\pi r^3}{3} \left[1 - \cos(\phi)\right]$$
$$V_{\text{outflow}} = 2.7 \times 10^{47} \text{m}^3 \left[0 \text{g}\right]$$
$$U_{\text{outflow}} \sim 1.5 \times 10^{-7} \text{ J m}^{-3} \left[1 - \cos(\phi)\right]$$

However, knots are ballistic Herbig-Haro objects. Total energy in knots $\sim 10^{37}$ J (Allen & Burton 1993). The large majority of the outflow energy is in the highly-collimated central outflow.



It seems that in OMC 1, if it doesn't have $\sim 10^{40}$ J of energy and an energy density of $\sim 10^{-7}$ Jm⁻³, it's not worth talking about:

$$U_{_B} \sim 1.6 \times 10^{-7} \text{ J m}^{-3}$$

 $U_{_G} \sim (0.5 - 8.8) \times 10^{-7} \text{ J m}^{-3}$
 $U_{_{outflow}} \sim 1.5 \times 10^{-7} \text{ J m}^{-3}$ (not uniformly distributed)

So how can we tell which forces determine the magnetic field morphology?

Alfvén and ballistic velocities $c_A = \frac{B}{\sqrt{\mu_0 \rho}}$ $B = 6.4 \pm 2.1 \text{ mG},$ $c_A = 9.1 \pm 3.0 \text{ km s}^{-1}$

Max. Alfvénic field deviation in 500 yr: (4.6 \pm 1.5) \times 10⁻³ pc

However, ejecta velocities are supersonic and super-Alfvénic: Typical LOS ejecta velocity is \sim 150 km/s (Bally et al. 2017).

Max. distance travelled in 500 yr: $\sim 10^2$ pc

Hence, the outflow **cannot** have caused the observed deviation in magnetic field lines in OMC 1 in the time since its formation.

Our results suggest that:

- The magnetic field strength and the gravitational force between BN/KL and S are in approximate balance
- The 'hourglass' magnetic field shape may have been produced by the gravitational interaction. The magnetic field may have been compressed until equilibrium was reached
- The orientation of the BN/KL outflow has, on large scales, been determined by the magnetic field



Finally: something about ALMA!

"Does magnetic field of the natal clump gas regulate outflows in a forming star cluster?"

- PI: Ray Furuya, Co-Is: Pattle, Hasegawa et al.
 - Band 7, 0.4-arcsec resolution dust continuum polarimetric imaging
 - 4 outflow-driving sources in OMC1 South
 - Grade B-ranked in Cycle 5



Conclusions:

- We measure a magnetic field strength of 6.6±4.7 mG in the OMC 1 region of Orion A, comparable to more distant high-mass star-forming regions
- We find that the magnetic, gravitational and outflow energy densities in OMC 1 are all $\sim 10^{-7}$ J m⁻³
- The Alfvén velocity in OMC 1 is sufficiently small that the deviations in the magnetic field of OMC 1 could not have been created in the lifetime of the BN/KL outflow
- We suggest that the 'hourglass' morphology of the magnetic field in OMC 1 is caused by distortion of an initially-uniform magnetic field by the gravitational interaction of Orion BN/KL and AS
- We further suggest that the orientation of the large-scale BN/KL outflow is constrained by the orientation of the magnetic field.
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For more detail: Pattle et al. 2017, ApJ 846 122

Thank you!



Velocity field in the OMC 1 region



HARP ¹³CO measurements; Buckle et al. 2010

Figure 13. (a) 13 CO intensity weighted velocity map of Orion A, giving the velocity field in km s⁻¹. The black line marks the region between

Velocity field in the OMC 1 region



HARP ¹³CO measurements; Buckle et al. 2010

There is more than one interpretation of this line-of-sight velocity...



Hence, we cannot exclude either hypothesis using velocity data alone.

GPE density: point-source model

From our column density map: Mass of BN/KL: 1000 M_{\odot} Mass of AS: 285 M_{\odot}

GPE of two point sources:
$$E_G =$$

Plane-of-sky separation = 0.166 pc; for a conservative estimate of filament orientation,
$$r = 0.166 \times (2)^{0.5} \sim 0.23$$
 pc

$$E_{G} = -1.0 \times 10^{40} J$$

 GM_1M_2

r

Assuming OMC 1 is a cylinder with height $0.321 \times (2)^{0.5}$ pc and diameter 0.141 pc:

 $U_{_{G}} = -0.5 \times 10^{-7} \text{ J m}^{-3}$

BN/KL

r = 0.23 pc

AS

GPE density: prolate spheroid model

Prolate-spheroid model:

$$E_G = -\frac{8}{15}\pi^2 G\rho^2 a_1^4 a_3 \times \frac{1}{e} \ln\left(\frac{1+e}{1-e}\right)$$

(e.g. Binney & Tremaine 2008)

Assuming the BN/KL-AS system is a prolate spheroid with mass 1465 M_{\odot} , semi-major axis $a_3 = 0.16$ pc and semiminor axes $a_1 = 0.071$ pc:

$$E_{G} = -8.6 \times 10^{40} \text{ J}$$

 $U_{G} = -8.8 \times 10^{-7} \text{ J m}^{-3}$

